

Part I:

Graphical Symplectic Algebra

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Motivation: I am bad at mathematics.

A **prop** is a strict symmetric monoidal category generated by a single object...



A **compact prop** also allows for wires to be bent/unbent:



Graphical linear algebra

Affine matrices: generators

Given a field \mathbb{K} , finite dimensional affine transformations can be represented their **homogeneous coordinates matrices** (T, S are matrices, \vec{a}, \vec{b} are vectors):

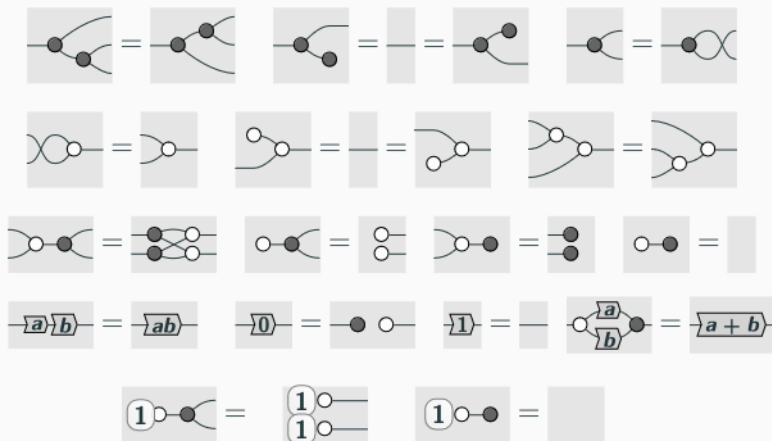
$$\left[\begin{array}{c|c} T & \vec{a} \\ \hline 0 & 1 \end{array} \right] \left[\begin{array}{c|c} S & \vec{b} \\ \hline 0 & 1 \end{array} \right] = \left[\begin{array}{c|c} TS & T\vec{b} + \vec{a} \\ \hline 0 & 1 \end{array} \right]$$

The prop of affine transformations between finite dimensional vector spaces is generated by the homogeneous coordinate matrices:

$$\begin{aligned} \left[\begin{array}{c|c} \bullet & \text{---} \\ \hline 1 & 0 \\ 1 & 0 \\ \hline & 1 \end{array} \right] &= \left[\begin{array}{c|c} 1 & 0 \\ \hline 1 & 0 \\ \hline & 1 \end{array} \right] & \left[\begin{array}{c|c} \circ & \text{---} \\ \hline 1 & 1 \\ \hline & 1 \end{array} \right] &= \left[\begin{array}{c|c} 1 & 1 \\ \hline & 1 \end{array} \right] & \left[\begin{array}{c|c} \text{---} & \vec{a} \\ \hline & 1 \end{array} \right] &= \left[\begin{array}{c|c} a & 0 \\ \hline & 1 \end{array} \right] \\ \left[\begin{array}{c|c} \bullet & \text{---} \\ \hline * & * \\ \hline & 0 \end{array} \right] &= \left[\begin{array}{c|c} * & * \\ \hline & 0 \end{array} \right] & \left[\begin{array}{c|c} \circ & \text{---} \\ \hline * & 0 \\ \hline & 0 \end{array} \right] &= \left[\begin{array}{c|c} * & 0 \\ \hline & 0 \end{array} \right] & \left[\begin{array}{c|c} \textcircled{1} & \circ \text{---} \\ \hline * & 0 \\ \hline & 1 \end{array} \right] &= \left[\begin{array}{c|c} * & 0 \\ \hline & 1 \end{array} \right] \end{aligned}$$

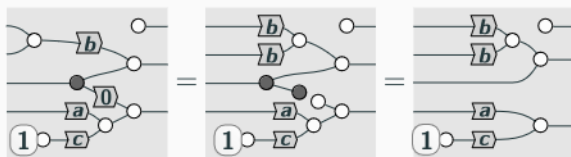
Affine matrices: axioms

Modulo the equations:



Example of matrix multiplication

The following diagram can be simplified to a normal form:

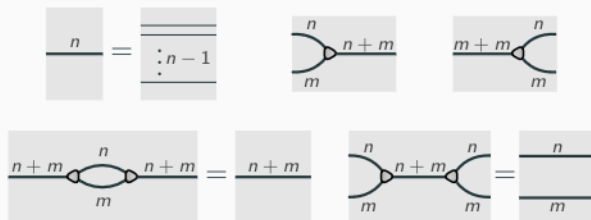


Following the paths from left to right gives us the homogeneous coordinate matrix:

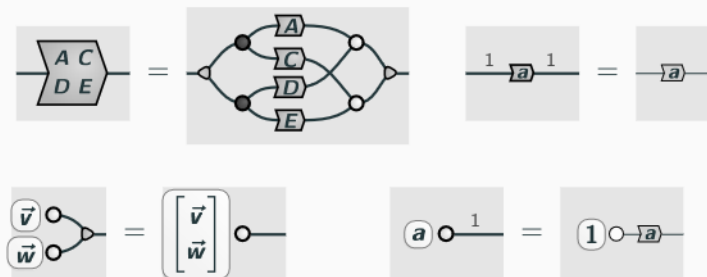
$$\begin{array}{c} x_0 \\ x_1 \\ x_2 \\ x_3 \\ \textcircled{1} \end{array} \begin{array}{c} \text{b} \\ \text{b} \\ \text{a} \\ \text{c} \end{array} \begin{array}{c} x_0 b \\ x_1 b \\ x_3 a \\ c \end{array} \begin{array}{c} 0 \\ x_0 b + x_1 b + x_2 \\ x_3 a + c \end{array} \rightsquigarrow \left[\begin{array}{cccc|c} 0 & 0 & 0 & 0 & 0 \\ b & b & 1 & 0 & 0 \\ 0 & 0 & 0 & a & c \\ \hline & & & & 1 \end{array} \right] \begin{array}{c} x_0 \\ x_1 \\ x_2 \\ x_3 \\ \hline 1 \end{array} = \begin{array}{c} 0 \\ x_0 b + x_1 b + x_2 \\ \hline x_3 a + c \\ 1 \end{array}$$

Strictification and block matrices

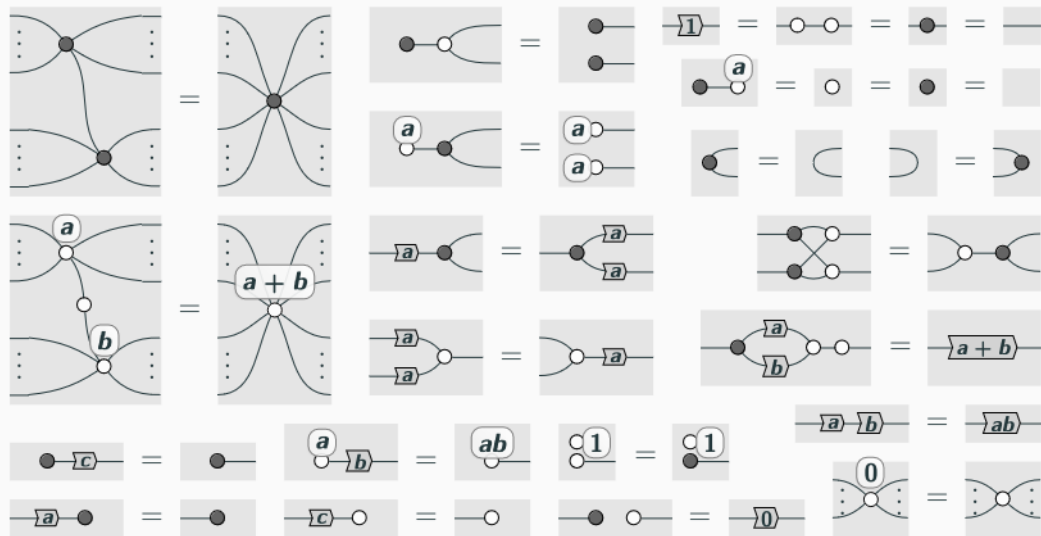
Every prop can be strictified to an \mathbb{N} -coloured prop:



This allows us to define block matrices/vectors diagrammatically:

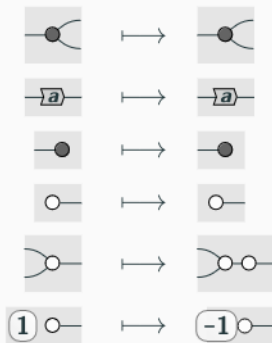


Modulo, the “spiders” $m \begin{array}{c} \circ \\ \text{---} \end{array} \begin{array}{c} \text{a} \\ \text{---} \end{array} \begin{array}{c} \circ \\ \text{---} \end{array} n$ and $m \begin{array}{c} \bullet \\ \text{---} \end{array} \begin{array}{c} \text{---} \\ \text{---} \end{array} \begin{array}{c} \bullet \\ \text{---} \end{array} n$ being commutative, undirected and,



for all $a, b \in \mathbb{K}$, $c \in \mathbb{K}^\times$.

The embedding $\text{AffMat}_{\mathbb{K}} \hookrightarrow \text{AffRel}_{\mathbb{K}}$ taking an affine transformation $T : n \rightarrow m$ to its graph $\{(\vec{x}, T\vec{x}) \mid \vec{x} \in \mathbb{K}^n\}$ sends:



Classical mechanics and symplectic geometry

The extensional behaviour of an electrical circuits is characterised by how it transforms current and voltage;

- **Ohm's law:** The voltage around the node in a circuit is equal to the current multiplied by the resistance.
- **Kirchhoff's current law:** The sum of currents flowing into a node is equal to the sum of currents flowing out of the node.

Example

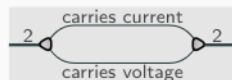
Given a linear resistor with resistance $r \in \mathbb{R}^{>0}$ on a wire with incoming current/voltage (z_0, x_0) and outgoing current/voltage (z_1, x_1) :

- by KCL, currents equalize: $z_0 = z_1$;
- by OL, the outgoing current becomes: $x_1 = x_0 + z_0 r$.

String diagrams for electrical circuits, take I

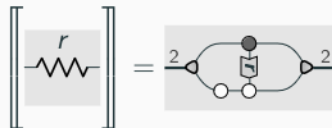
Following Baez et al. [BCR18] and Baez and Fong [BF18], we can represent electrical circuit components as real affine relations.

Using the string diagrams from Bonchi et al. [Bon+19], decompose a wire into a current and voltage



...the resistor is represented as follows:

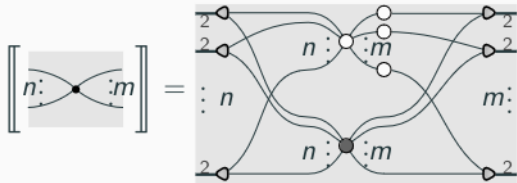
Example



More string diagrams for electrical circuits, take I

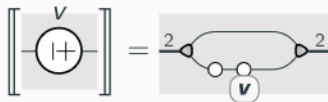
Example

Ideal wire junctions *sum currents*, and *equalize voltages*:



Example

Constant voltage source *does nothing to current* and *adds to the voltage*:



What is the more conceptual picture?

Classical mechanical systems can be represented by the configurations of abstract **positions** Z and **momenta** X :

<i>Classical mechanics</i>	Z	dZ/dt	X	dX/dt
Translation	position	velocity	momentum	force
Electronic	charge	current	flux linkage	voltage
Hydraulic	volume	flow	pressure momentum	pressure
Thermal	entropy	entropy flow	temperature momentum	temperature

For n -particles in Euclidean space, *the space of possible configurations of positions/momenta* $\mathbb{R}^{2n} \cong \mathbb{R}_Z^n \oplus \mathbb{R}_X^n$ is the **phase space**.

Table adapted from Smith [Smi93, page 23, table 2.1] and Baez and Fong [BF18]

Definition

Two configurations $(\vec{z}, \vec{x}), (\vec{q}, \vec{p}) \in \mathbb{K}^{2n}$ of phase-space are **compatible** when:

$$\vec{z} \cdot \vec{p} - \vec{x} \cdot \vec{q} = 0$$

The bilinear map

$$\omega_n : \mathbb{K}^{2n} \oplus \mathbb{K}^{2n} \rightarrow \mathbb{K} \quad ((\vec{z}, \vec{x}), (\vec{q}, \vec{p})) \mapsto \vec{z} \cdot \vec{p} - \vec{x} \cdot \vec{q}$$

is a **symplectic form**, and the phase space $(\mathbb{K}^{2n}, \omega_n)$ is a **symplectic vector space**.

An **affine Lagrangian subspace** is a *maximally compatible* affine subspace of a symplectic vector space.

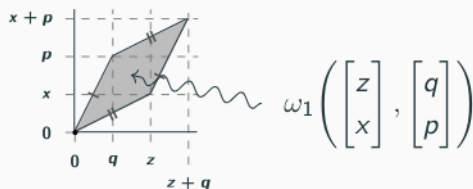
Remark (Baez and Fong [BF18], Baez et al. [BCR18])

Resistors, voltages sources and junctions of wires are affine Lagrangian subspaces.

Geometric interpretation of compatibility

Example

In the phase-space of a single particle, (\mathbb{K}^2, ω_1) , the symplectic form measures area:



Compatible points are colinear, so affine Lagrangian subspaces are lines.

An affine Lagrangian subspaces don't represent single particle; but an ensemble of particles *flowing along a trajectory*.

Definition (Guillemin and Sternberg [GS79], Weinstein [Wei82])

The compact prop of affine Lagrangian relations $\text{AffLagRel}_{\mathbb{K}}$ has:

- **Morphisms** $n \rightarrow m$, given by (possibly empty) affine Lagrangian subspaces of $(\mathbb{K}^{2n} \oplus \mathbb{K}^{2m}, \omega_n - \omega_m : \mathbb{K}^{2(n+m)} \oplus \mathbb{K}^{2(n+m)} \rightarrow \mathbb{K})$.
- **Composition** is given by relational composition.
- **Symmetric monoidal structure** is given by the direct sum.

Lemma

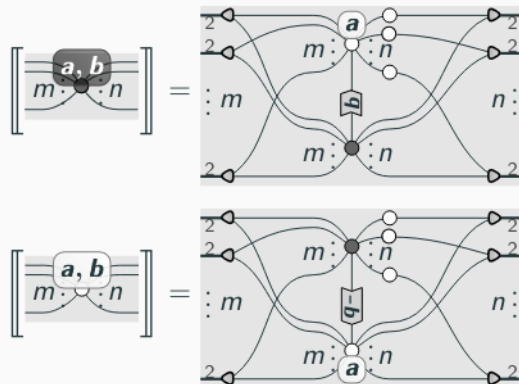
There is an embedding $\text{AffRel}_{\mathbb{K}} \rightarrow \text{AffLagRel}_{\mathbb{K}}$ given

- **on objects by:** $n \mapsto 2n$;
- **on morphisms by:** $(S + \vec{a}) \mapsto S^{\perp} \oplus (S + \vec{a})$.

For the geometrically inclined, this is induced by the embedding of a vector space $\mathbb{R}^n \hookrightarrow T^(\mathbb{R}^n) \cong (\mathbb{R}^n)^* \oplus \mathbb{R}^n \cong \mathbb{R}^{2n}$ into its cotangent bundle.*

Generators of affine Lagrangian relations (Comfort and Kissinger [CK22])

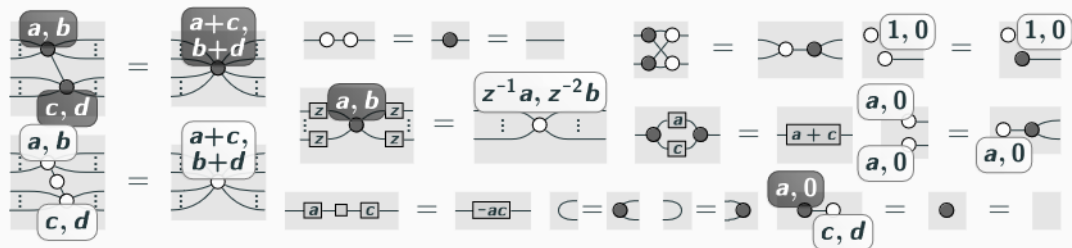
$\text{AffLagRel}_{\mathbb{K}}$ is generated by two spiders decorated by \mathbb{K}^2 ; interpreted in $\text{AffRel}_{\mathbb{K}}$ as:



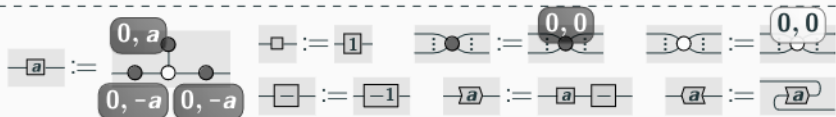
Equations of affine Lagrangian relations (Booth et al. [BCC24b])

Modulo both spiders, being commutative, undirected nodes,

as well as for all $a, b, c, d \in \mathbb{K}$ and $z \in \mathbb{K}^\times$:

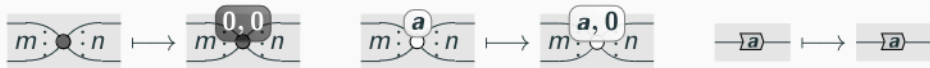


With derived
generators:

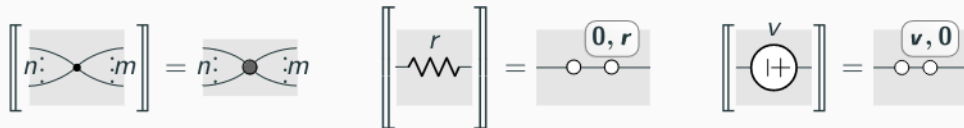


Interpreting electrical circuits

The embedding $\text{AffRel}_{\mathbb{K}} \hookrightarrow \text{AffLagRel}_{\mathbb{K}}$ takes:



Now that the position/momentum wires are bundled together, we have a more concise description of electrical circuit components:

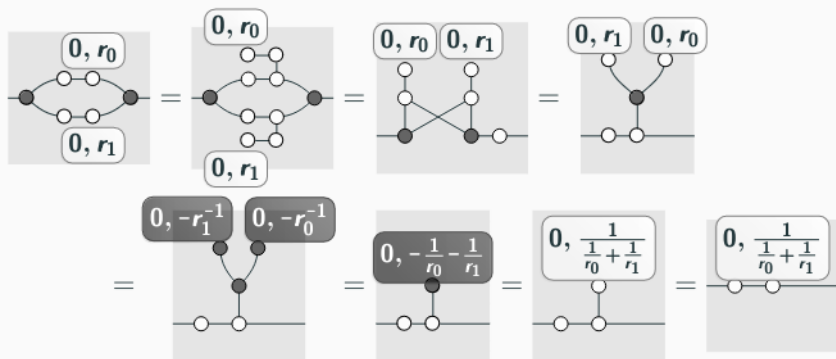


Example: composing resistors in parallel

AffLagRel $_{\mathbb{R}}$ allows us to cleanly compose electrical circuits:

Example

Consider two resistors with resistances $r_0, r_1 \in \mathbb{R}^{>0}$ composed in parallel.



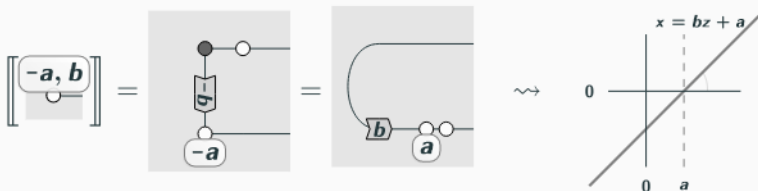
Electrons nondeterministically flow through both resistors, where they are impeded.

They extensionally behave like a resistor with resistance $1/(1/r_0 + 1/r_1)$.

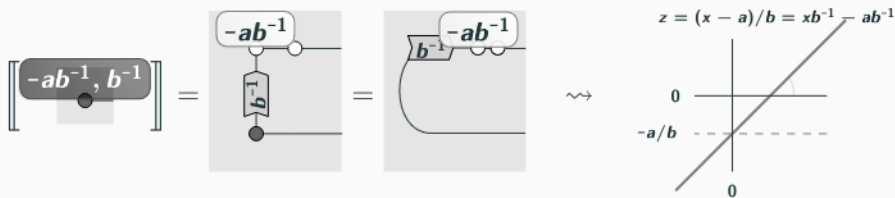
Interpreting the two spiders:

This colour-swap rule corresponds to a change of reference frame.

Where configurations of phase space can be represented as functions of position:



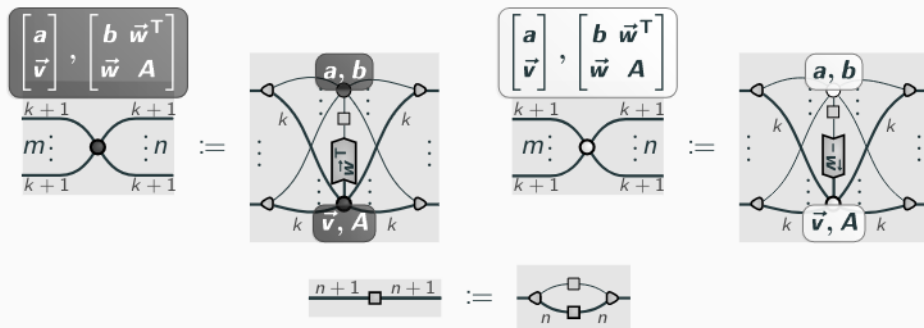
...or of momentum:



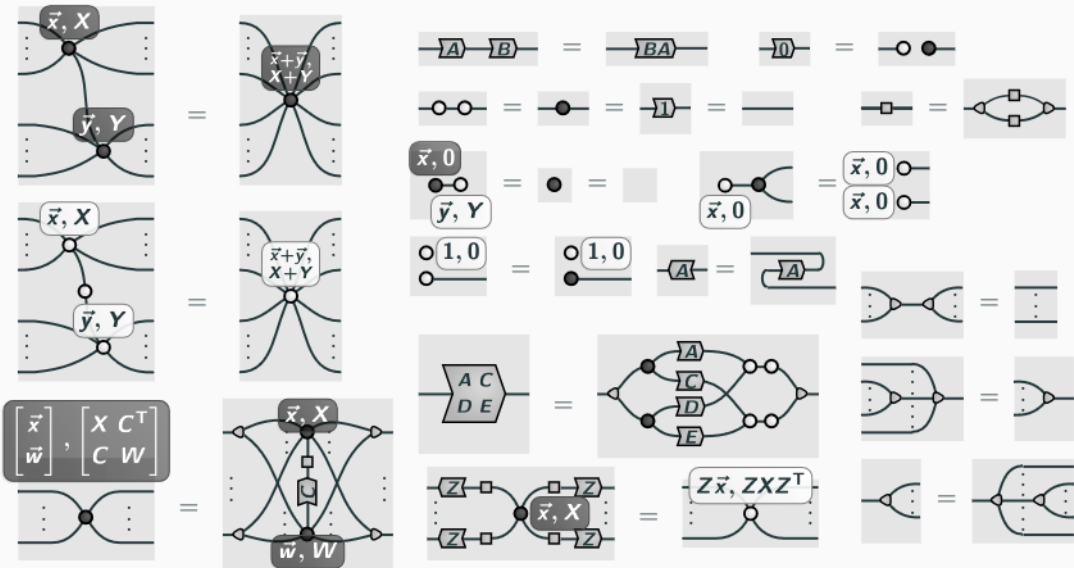
Scalable spiders

We can define higher-dimensional spiders by induction on the number of wires $k \in \mathbb{N}$.

Take $n, m \in \mathbb{N}$, $a, b \in \mathbb{K}$, $\vec{v}, \vec{w} \in \mathbb{K}^k$ and $A \in \text{Sym}_k(\mathbb{K})$.



Scalable identities



Impedance matrix

Consider a network of resistors/voltage sources acting on n wires.

The extensional behaviour can be represented by a positive-definite $0 \prec R \in \text{Sym}_n(\mathbb{R})$ called the **impedance matrix**, and a voltage $\vec{v} \in (\mathbb{R}^{>0})^n$

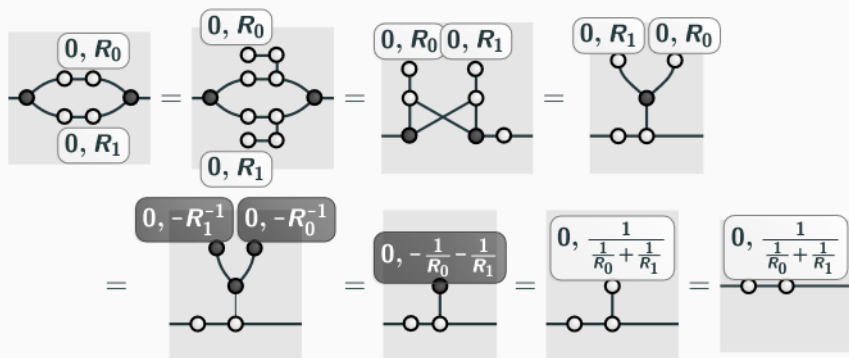
$$\left[\begin{array}{c} \vec{v}, R \\ \text{---} \circ \text{---} \circ \text{---} \\ n \qquad n \end{array} \right] = \left\{ \left(\begin{bmatrix} \vec{z} \\ \vec{x} \end{bmatrix}, \begin{bmatrix} \vec{z} \\ \vec{x} + R\vec{z} + \vec{v} \end{bmatrix} \right) \mid \forall \vec{z}, \vec{x} \in \mathbb{R}^n \right\}$$

The resistance between the j th and k th wire is $r_{j,k} = r_{k,j} \in \mathbb{R}$.

The change in voltage on wire j is $v_j \in \mathbb{R}$.

Composing networks of resistors in parallel

Black-boxed networks of resistors compose in parallel in the same way as single resistors composed in parallel:



We don't know the internal structure of the two networks, but we still can compute their extensional behaviour in parallel.